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A new multifractal-based Grain Size Distribution model

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33 Abstract

Previous works related to the application of the multifractal theory for analyzing the grain size 34 distribution (GSD), showed the potential of this approach to deal with this complex issue. 35 However, absence of the practical application of this kind of statistical analysis raised some 36 37 doubts among the soil scientists. Compared to the experimental dry sieving method, which is 38 based on mass representations of different grain sizes, the approach presented in this work relies on the analysis of grain densities (density indicators) scanned by means of X-ray CT 39 (Computed Tomography). By reducing the resolution of the scanned soil image(s), the 40 41 cumulative representation of solid particles equal to or larger than the actual discretization element can be determined, and described analytically by means of the universal multifractals 42 (UM). 43

For validation of the new UM approach, the X-ray CT results of three different soils were 44 used: the volcanic substrate covering Green Wave (a green roof of Champs-sur-Marne in 45 46 France), and two horizons of the soil collected from the low land mountain area of Sierra de 47 Guadarrama in Spain. Comparison between the proposed UM model and the experimental data of these three materials confirms that the GSD can be reasonably well predicted from the 48 scanned images of soils covering wide range of grain sizes. The UM model, unlike the fractal-49 based models, accounts for fractal dimension that depends on grain size, and hence, based on 50 the preliminary results presented in this work, it could be rather useful in case of multi-modal 51 52 soils whose GSD curves are described with multiple fractal dimensions.

53

54 **Key words:** grain size distribution; multifractals; X-ray computed tomography; granular

55 media

56 1. Introduction

The grain size distribution (GSD) is one of the fundamental properties of granular soils that, 57 58 besides the influence on mechanical characteristics, also affects the packing arrangement of grains (Nolan & Kavanagh, 1993; He et al., 1999 among the others), and hence the 59 60 distribution of pores that further impacts the hydraulic properties of the porous medium (Segal et al., 2009). Most often the GSD curve is experimentally determined based on the mass 61 62 fractions of different grain sizes extracted either by using sieves of different void sizes, for 63 grains larger than 80 microns (dry sieving method - AFNOR, 1996), or by means of 64 sedimentation test (AFNOR, 1992; Beuselinck et al., 1998) for finer particles. The alternative 65 approach proposed to measure GSD is a laser diffraction method (Miller & Schaetzl, 2012).

Detailed overview of different approaches used for describing the complexity of GSD curves 66 67 can be found in Ghanbarian & Hunt (2017). One of them is the self-similarity principle which is included in fractal-based models and which assumes occurrence of the same pattern of the 68 soil structure at all scales. According to Ghanbarian-Alavijeh et al. (2011), the three-phase 69 PSF (pore-solid-fractal) approach (Perrier et al., 1999; Bird et al., 2000) is the most consistent 70 71 and with the strongest physical-basis among the fractal-based approaches. Besides pores and grains, it assumes one additional "fictive" type of soil elements - fractals - that are 72 73 successively broken at smaller scales in a self-similar way, leading finally to the structure 74 consisting of fractal-distributed pore and grain sizes. Thus, the GSD can be represented by means of a power (fractal) law, where the fractal coefficient is included in the exponent. 75 76 However, unlike assumed in the PSF model, grain densities are non-homogeneous, which also contributes to the complexity of distribution of different mass fractions that often cannot be 77 described with a single fractal dimension (Bittelli et al., 1999). 78

79 Multifractal formalism, that takes into account different fractal coefficients for different 80 threshold values, was also used for analyzing the complexity of GSD. Grout et al. (1998) and Posadas et al. (2001) used Renyi dimensions, one of the multifractal parameters, to 81 82 characterize the heterogeneous distribution of different mass fractions. Besides this type of multifractal analysis, the singularity spectra analysis is also applied for analyzing the dry soil 83 volume-size distribution obtained by using a laser distraction method (Martín & Montero, 84 2002). Recently, Torre et al. (2016) used a X-ray CT, a non-destructive technique for 85 86 obtaining a three-dimensional grey-scale image of a porous material (Hseih, 2003; Banhart, 2008) in order to compare the three-dimensional structural complexity of spatial arrangement 87 88 of grains and pores, with that of differently oriented two-dimensional planes. The multifractal analysis has also proved to be convenient in this case. Even though the multifractal theory 89 brings great potential for understanding the complexity of GSD (Ghanbarian & Hunt, 2017), 90 91 up to date this kind of analysis has not found practical application.

92 This work is focused on development of a new physically-based GSD model founded on the 93 Universal Multifractal (UM) framework (Schertzer & Lovejoy, 1987; Schertzer & Lovejoy, 1997). Based on a grey-scale soil image scanned by means of X-ray CT, it is possible to 94 95 recognize solid particles of different sizes by progressively decreasing the resolution of the image while keeping the fixed value of the threshold. Change of the representation of solid 96 97 particles with the resolution of the image can be directly linked with the grain size 98 distribution, and described analytically in a mathematically-elegant way by means of the UM 99 framework. Compared to work of Lai & Chen (2018), where a sophisticated machine learning 100 tool was used for particle recognition, this approach is much simpler and more convenient for 101 practical application.

102 The UM framework in combination with X-ray CT imaging was firstly validated for some
103 artificial volcanic substrate (Stanić et al., 2019; Stanić et al., 2020) used for covering green

roof named Green Wave (Versini et al., 2018; Versini et al., 2020). Results of the model, 104 105 whose parameters are directly determined from scanned images, were first compared with the 106 experimental data obtained by means of the standard dry sieving method (AFNOR, 1996) and 107 sedimentation test (AFNOR, 1992). Furthermore, the UM model was tested on scanned 108 images of two horizons of an intact soil sample collected from the low land mountain area of 109 Sierra de Guadarrama (Schmid et al. 2016) called La Herreria. In this case, results of the model were compared with measured percentages of sand, silt and clay particles, since 110 111 detailed GSD curves are lacking. Finally, for published experimental GSD data of the GW substrate and Walla Walla soil (Bittelli et al., 1999), the UM model was compared with the 112 113 fractal-based PSF model.

114 **2.** Methodology

The GSD model proposed in this work is based on the recognition of solid particles of 115 different sizes from the scanned soil image, by changing the resolution of the image. This can 116 be described analytically through the application of the Universal Multifractals (UM) 117 118 framework (Schertzer & Lovejoy, 1987; Schertzer & Lovejoy, 1997) which is briefly 119 described below. Note that two-dimensional scanned soil images (Euclidian dimension E = 2), 120 extracted from the three-dimensional one (E = 3), were analyzed in this work in order to simplify the methodology presented. However, this simplified approach (E = 2) is credible 121 122 only under certain conditions that are described later in the text, while otherwise the same 123 methodology should be applied for E = 3.

124 2.1. Universal Multifractal (UM) theoretical framework

125 In Figure 1 is presented a renormalized two-dimensional soil density indicator field $\rho^{ind}(\lambda)$ at 126 various resolutions λ , which is, due to the better visualization, presented in a three-

dimensional form. Here, $\lambda = \frac{L}{l(\lambda)}$ is equal to the ratio between the size of the image L [L] and 127 128 the size of a single pixel $l(\lambda)$ [L], representing the number of pixels along an edge of the image. Values of $\rho^{ind}(\lambda) = \rho(\lambda) / \rho_{bulk}$ are presented as histogram at each λ , where ρ_{bulk} is the 129 constant bulk density of the dry material $[M/L^3]$ ($\rho^{ind} = 1$ is a renormalized ρ_{bulk}). Clearly, 130 $\rho^{ind}(\lambda)$ values mitigate as λ decreases by merging pixels in groups by λ_1^2 , where λ_1 is an 131 integer value equal 2 (check dashed gridlines in Figure 1E). By averaging $\rho^{ind}(\lambda)$ values of 132 each group, attenuated $\rho^{ind}(\lambda)$ field is obtained with λ_l times smaller λ and λ_l times larger 133 pixel size $l(\lambda)$. 134

By means of UM (Schertzer & Lovejoy, 1987; Schertzer & Lovejoy, 1997) it is possible to compute, at different λ , the probability that $\rho^{ind}(\lambda)$ exceeds λ^{γ} , a renormalized threshold value that changes with λ (a transparent color platform in Figure 1). Threshold value is expressed through the fixed dimensionless singularity γ [-], and therefore it decreases together with λ (see Figure 1a to Figure 1f) until it reaches unit value at $\lambda = 1$. For a certain value of γ (= 0.211 in Figure 1), the previously mentioned probability of exceeding can be computed as:

141
$$P(\rho^{ind}(\lambda) \ge \lambda^{\gamma}) = \frac{N(\rho^{ind}(\lambda) \ge \lambda^{\gamma})}{\lambda^{E}} \approx \lambda^{-c(\gamma)}$$
(1)

142
$$c(\gamma) = C_1 \left(\frac{\gamma}{C_1 \alpha'} + \frac{1}{\alpha}\right)^{\alpha'}; \quad \alpha' = \left(1 - \frac{1}{\alpha}\right)^{-1}, \quad \alpha \neq 1$$
 (2)

143 where $N(\rho^{ind}(\lambda) \ge \lambda^{\gamma})$ is the number of $\rho^{ind}(\lambda)$ values that are equal to or higher than λ^{γ} , 144 while $c(\gamma)$ is the co-dimension function that, besides γ , depends on two parameters, C_I and α . 145 Note that in Equation (1) an approximate equal sign is used because term $\frac{N(\rho^{ind}(\lambda) \ge \lambda^{\gamma})}{\lambda^{E}}$ is 146 computed by counting $N(\rho^{ind}(\lambda) \ge \lambda^{\gamma})$ at different λ (discrete form), while $\lambda^{-c(\gamma)}$ is related 147 to the UM analytical form.

Parameters C_1 and α fully characterize ρ^{ind} field, where C_1 describes the sparseness of the 148 mean value of the field while α describes the change of sparseness for values around the 149 150 mean. As explained in Schertzer & Lovejoy (1987), C_1 takes values between 0 (mean value is 151 ubiquitous - homogeneous field) and E (mean value is too sparse to be observed), while α takes values between 0 (no occurrence of extremes - fractal field) and 2 (maximal occurrence 152 of extremes – log-normal field). Equation (1) is presented in Figure 2a for $C_1 = 2.23 \times 10^{-2}$ and 153 $\alpha = 1.67$, values that characterize $\rho^{ind}(\lambda)$ field in Figure 1, and different γ values (including $\gamma =$ 154 0.211) corresponding to various dashed lines. 155

156 2.2. Adaptation of the UM framework – New GSD model

157 In this work, the presented up-scaling procedure is used for recognizing solid particles of different sizes from the obtained $\rho^{ind}(\lambda)$ field. Compared to the previous explanation, where a 158 resolution dependent threshold λ^{γ} was accounted for, here is used a fixed threshold value 159 $\rho_{s,min}^{ind}$ (solid platform in Figure 1) related to the renormalized minimal grain density ($\rho_{s,min}^{ind} >$ 160 1). Therefore, $\rho^{ind} \ge \rho_{s,min}^{ind}$ values indicate the total area of the image covered with solid 161 particles (grains). As shown in Figure 1, this area reduces when up-scaling, mostly by getting 162 rid of isolated $\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind}$ values that are surrounded by those lower than $\rho_{s,min}^{ind}$. On the 163 contrary, larger continuous zones covered by $\rho^{ind}(\lambda) \ge \rho^{ind}_{s,min}$ values resist longer to the up-164 165 scaling process, indicating the presence of a large grain on that location (central zone in Figure 1). Therefore, the total area covered with $\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind}$ values at certain λ indicates 166 a cumulative representation of solid particles of diameter equal to or larger than the size of a 167 168 single pixel $l(\lambda) = L / \lambda$:

169
$$P(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind}) = \frac{N(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind})}{\lambda^{E}}$$
(3)

170 In order to transform Equation (3) into the distribution function $P (d \ge l(\lambda))$, it is necessary to 171 renormalize it with respect to the initial representation of solid particles met at $\lambda_n \le \lambda_{up}$. 172 Therefore, $P (d \le l(\lambda)) = 1 - P (d \ge l(\lambda))$ can be expressed as:

173
$$P(d < l(\lambda)) = 1 - \frac{P(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind})}{P(\rho^{ind}(\lambda_n) \ge \rho_{s,min}^{ind})} = 1 - \frac{N(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind})}{N(\rho^{ind}(\lambda_n) \ge \rho_{s,min}^{ind})} \left(\frac{\lambda_n}{\lambda}\right)^E$$
(4)

174 The analogy between Equation (4) and the dry sieving method is explained in the Appendix. 175 Equation (3), and hence Equation (4), can be described analytically through the UM 176 framework if expressing $\rho_{s,min}^{ind}$, which is independent on λ , through λ :

177
$$\rho_{s,min}^{ind} = \frac{\rho_{s,min}}{\rho_{bulk}} = \lambda^{\gamma(\lambda)}$$
(5)

where $\rho_{s,min}$ is the minimal grain density [M/L³], and $\gamma(\lambda)$ differs from a fixed γ used in Equations (1) and (2), since it changes with λ in order to maintain fixed value of $\rho_{s,min}^{ind}$:

180
$$\gamma(\lambda) = \frac{ln(\rho_{s,min}^{ind})}{ln(\lambda)}$$
 (6)

181 Finally, by introducing Equation (6), instead of γ, into Equation (2), it is possible to express
182 Equation (3) analytically:

183
$$P(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind}) \approx \lambda^{-c(\gamma(\lambda))} = \lambda^{-C_1 \left(\frac{\ln(\rho_{s,min}^{ind})}{C_1 \alpha'} + \frac{1}{\alpha}\right)^{\alpha'}}$$
(7)

Equation (7) is presented in Figure 2b with solid line which is also fully characterized by means of parameters C_1 and α . Value of $\gamma(\lambda) = C_1$ corresponds to the upper resolution limit $\lambda_{up} = (\rho_{s,min}^{ind})^{1/C_1}$ (see Equation 6) for which, due to the fact that $c(\gamma(\lambda) = C_1) = C_1$ (see Equation 2), $P(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind})$ reaches its maximal value equal to $(\rho_{s,min}^{ind})^{-1}$. On the 188 contrary, the lower resolution limit $\lambda_{low} = (\rho_{s,min}^{ind})^{1/\gamma_s}$ can be also computed from Equation 189 (6) for $\gamma(\lambda) = \gamma_s$, which is known as the most probable singularity:

190
$$\gamma_s = C_1 \alpha' \left(\left(\frac{E}{C_1} \right)^{1/\alpha'} - \frac{1}{\alpha} \right)$$
 (8)

- 191 Note that Equation (8) is derived from Equation (2) given that $c(\gamma_s) = E$.
- Finally, having on mind that $l(\lambda) = d_g$, and thus $\lambda = L / d_g$, the analytical GSD function can be derived by introducing Equation (7) into (4):

$$194 \qquad P(d < d_g) = 1 - \frac{\left(\frac{ln(\rho_{s,min})}{ln(L/d_g)} + \frac{1}{\alpha}\right)^{\alpha'}}{\left(\frac{L}{d_g,min}\right)^{-C_1} \left(\frac{\frac{ln(\rho_{s,min})}{ln(L/d_{g,min})}}{c_1\alpha'} + \frac{1}{\alpha}\right)^{\alpha'}}$$
(9)

195 where $d_{g,min} = l(\lambda_n) = L / \lambda_n$ is the minimal grain diameter [L] equal to the size of a pixel at λ_n . 196 From Equation (9) the probability density function can be derived as the first derivative of 197 $P(d \ge d_g) = 1 - P(d < d_g)$ with respect to $ln(L/d_g)$, providing the following expression:

$$198 \qquad p(d = d_g) = -c \left(\frac{L}{d_g}\right) \left(\frac{L}{d_{g,min}}\right)^{c \left(\frac{L}{d_{g,min}}\right)} \left[1 - \frac{\frac{ln(\rho_{s,min}^{ind})}{C_1}}{\frac{1}{\alpha} ln\left(\frac{L}{d_g}\right) + \frac{ln(\rho_{s,min}^{ind})}{C_1\alpha'}}\right] e^{-c\left(\frac{L}{d_g}\right) ln\left(\frac{L}{d_g}\right)} \tag{10}$$

199 where
$$c\left(\frac{L}{d_g}\right) = C_1 \left(\frac{ln(\rho_{s,min}^{ind})}{ln(L/d_g)C_1\alpha'} + \frac{1}{\alpha}\right)^{\alpha'}$$
 is the co-dimension, while $E - c\left(\frac{L}{d_g}\right)$ describes the

change of fractal dimension with d_g . The approach proposed here would face certain issues mostly related to the way pixels are grouped. Therefore, it is possible to have λ_l^2 neighbor pixels that belong to a grain of larger size, but since they are distributed in different groups there is a "good" chance that this larger grain will not be recognized after the aggregation. On 204 the contrary, those pixels can signify separated grains, but if they are aggregated as a part of 205 the same group of λ_I^2 pixels, they will be recognized as a part of the larger grain. 206 Nevertheless, these special cases do not influence the proposed algorithm significantly if 207 applied on the sufficiently large λ .

208 2.2.1. Determination of the model parameters

Parameters of the proposed model (Equation 9) are: $d_{g,min}$, $\rho_{s,min}^{ind}$, α and C_1 . The first two are physical parameters whose values are either estimated based on the type of material ($d_{g,min}$), or calculated based on the experimentally determined values of ρ_{bulk} and $\rho_{s,min}$ (see Equation 5), while the last two (α and C_1) are statistical parameters determined by analyzing the scaling behavior of ρ^{ind} field. This can be done by means of Trace Moment (TM) technique (Schertzer & Lovejoy, 1987) which assumes that the scaling of the average statistical moments of order $p \langle (\rho^{ind}(\lambda))^p \rangle$ can be described through the moment scaling function K(p):

216
$$\langle \left(\rho^{ind}(\lambda)\right)^p \rangle \approx \lambda^{K(p)}$$
 (11)

217
$$K(p) = \frac{c_1}{\alpha - 1} (p^{\alpha} - p); \quad \alpha \neq 1$$
 (12)

where K(p) is described through parameters C_1 and α (for more details see Schertzer & Lovejoy, 1987). Note that $c(\gamma)$ and K(p) functions are linked by Legendre transform (Frisch & Parisi, 1985), meaning that for each γ there is a corresponding p (i.e. for $\gamma = C_1$ and $\gamma = \gamma_s$ the corresponding values are p = 1 and $p = p_s$, respectively).

To determine values of α and C_1 for a certain $\rho^{ind}(\lambda)$ field, the field is firstly up-scaled as previously described, and all $\rho^{ind}(\lambda)$ values are raised on a power p at each λ . The average value of such a modified field $\langle (\rho^{ind}(\lambda))^p \rangle$ is computed at each λ , and the procedure is repeated for variety of $p \ge 0$ values. After plotting $log(\langle (\rho^{ind}(\lambda))^p \rangle)$ against $log(\lambda)$, 226 different linear regressions depending on *p* value are formed. Their slopes are related to *K*(*p*) 227 values that form the moment scaling function. Based on Equation (12), the first derivative of 228 the obtained *K*(*p*) function at *p* = 1 is equal to $C_1 = \frac{dK(p)}{dp}|_{p=1}$ (calculated numerically), while 229 the ratio between the second and the first derivative at *p* = 1 is $\alpha = \frac{1}{C_1} \frac{d^2 K(p)}{dp^2}|_{p=1}$.

230 2.2.2. Influence of the model parameters

To better understand the influence of the four parameters on the model behavior, Equation (9) has been tested on different values of each parameter, as illustrated in Figure 3. For all cases presented in Figure 3, value of L = 100 mm is kept constant while changing values of the four model parameters.

The impact of C_1 on the GSD is illustrated in Figure 3a by increasing (dash-dotted line) / 235 236 decreasing (dashed line) its initial value (solid line) by 50 % while preserving values of the 237 three remaining parameters. Similarly, in Figure 3b value of α is changed by 50 % in both ways. Figure 3a shows that parameter C_1 mostly affects the break onto the finer particles and 238 the shape of that part of the curve in a way that smaller C_1 secures higher contribution of fine 239 240 grains (dashed line), while the case is opposite for higher C_1 (dash-dotted line). On the 241 contrary, the change of parameter α (Figure 3b) is less affecting the representation of small 242 grains, but it is mainly responsible for the slope of the central part of the GSD curve, where smaller α provides steeper curve. Thus, in case of granular soils higher α and smaller C_1 243 values describe well-graded, while smaller α and higher C_1 describe more uniformly graded 244 materials. Indeed, well-graded materials usually have lower total porosity due to the better 245 spatial packing of grains, meaning the lower representation of zeros in ρ^{ind} field that causes 246 stronger variability of the field (higher α) and lower intermittency of its mean value (lower 247 248 C_{1}).

The impacts of $\rho_{s,min}^{ind}$ and $d_{g,min}$ on the GSD curve are also tested by varying one of the parameters while maintaining the rest. As illustrated in Figure 3c, the higher $\rho_{s,min}^{ind}$ (more strict threshold value), the higher values of P ($d < d_g$) (Equation 9), and vice versa. Unlike the three other parameters, $d_{g,min}$ dictates the total range of scales ($L / d_{g,min}$) by affecting mostly the distribution of small grains (tail of the GSD) - see Figure 3d.

3. Soil sampling and Image acquisition

In this section are given information about soil sampling and image acquisition for threedifferent materials: Green Wave substrate and Horizons A and A20 of La Herreria soil.

257 **3.1.** Green Wave substrate

258 Green Wave substrate is an artificial coarse material (VulkaTec Riebensahm GmbH 2016) 259 with 4 % of organic matter, used for covering green roofs. Due to its volcanic nature (values of grain and dry bulk densities are 2.35 Mg/m³ and 1.42 Mg/m³, respectively), this material 260 does not create a significant load on the roof construction which is the reason it has been used 261 in case of Green Wave (Versini et al., 2018; Versini et al., 2020), a wavy shape green roof 262 263 located next to Ecole des Ponts ParisTech in Champs-sur-Marne, France. The substrate 264 contains 50 % of grains larger than 1.6 mm, with 10 % of particles between 10 and 20 mm in the coarse range, and 13 % of fine particles smaller than 80 µm. Distribution of grains larger 265 266 than 80 µm was determined by means of the dry sieving method (AFNOR, 1996), while the 267 sedimentation test (AFNOR, 1992) was used for finer particles. The curvature and uniformity coefficients are $C_c = (D_{30})^2 / (D_{60} \times D_{10}) = 1.95$ and $C_u = D_{60} / D_{10} = 55$, respectively, and hence 268 this substrate is regarded as well graded according to the ASTM D2487-06 (2006) standard. 269

270 The sample of the GW substrate (10 cm diameter and 15 cm height) was prepared by mixing 271 and pouring up the material into the plexiglass cylinder (compacted to in situ value of ρ_{bulk} =

1.42 Mg/m³), simulating the way substrate is placed on the roof to avoid segregation of fine 272 273 particles at the bottom. Tomographic scans were conducted with a RX Solutions Ultratom 274 microtomograph, including a Hamamatsu L10801 X-ray source and a Paxscan Varian 2520V flat panel detector (1920 x 1560 pix², pixel size 127 μ m). X-ray source tension and current 275 276 were respectively 200kV and 280µA. The detector was set at 4 fps, each projection resulting 277 of an average over 25 projections, giving a total number of 4320 averaged projections. The sample being a long cylinder, stack mode was used and set to three turns. The reconstructed 278 279 3D image is finally represented by 1785x1785x3072 voxels with the edge length of 53.7 µm.

280 **3.2.** La Herreria soil (Horizons A and A20)

281 Two intact samples (60 mm diameter and 100 mm height) of La Herreria soil were collected 282 in the low land mountain area of Sierra de Guadarrama in Spain (Schmid et al. 2016)., which 283 is a highly degraded type of site because of the livestock keeping. One soil sample was 284 extracted from the top 18 cm layer (Horizon A), being the result of biological alteration with 285 roots resulting in fertile soil. This layer is moderately acid, with 2.5% of organic matter, 0.8% 286 of Fe₂O₃, sandy texture (65% sand, 25% silt, 10% clay) and bulk density of 1.6 Mg/m³. The 287 second soil sample was extracted from 18-40 cm depth (Horizon A20), also presenting an acid 288 character (pH = 6) with 0.5% of organic matter, 0.7% of Fe₂O₃, 55%, 30% and 15% of sand, 289 silt and clay particles, respectively, and bulk density of 1.7 Mg/m³.

X-ray CT scanning was performed using a Phoenix v | tome | x m 240 kV system (GE Sensing
& Inspection Technologies GmbH, Wunstorf, Germany) at the Hounsfield Facility, University
of Nottingham, UK. The scanner consisted of a 240kV microfocus X-ray tube fitted with a
tungsten reflection target and a DXR 250 digital detector array with 200 µm pixel size (GE
Sensing & Inspection Technologies GmbH, Wunstorf, Germany). A maximum X-ray energy
of 140kV and 200 µA was used to scan the soil core. A total of 2400 projection images were

acquired over a 360° rotation. Each projection was the average of six images acquired with a detector exposure time of 200 ms and the resulting isotropic voxel edge length was 32 μ m. The 3D image of the soil samples used in this work is represented by $676 \times 676 \times 300$ voxels.

299 4. Results and Discussion

The approach presented in this work is firstly validated on soil images of the GW substrate, and the experimental GSD data of the same material. Then, it is applied on two horizons of La Herreria soil, but in this case only measured percentages of sand, silt and clay particles were compared with model results because detailed GSD data are lacking. Finally, the comparison with the fractal-based PSF model (Perrier et al., 1999; Bird et al., 2000) is presented.

305 4.1. Obtaining ρ^{ind} field by means of X-ray CT

The ρ^{ind} field is transformed from a grey-scale image carrying the information about different intensities of grey color - bright shades of grey represent high, and dark shades low density zones. Since grey-level intensities (*GL*) can be linked with density ρ by means of the linear regression (Taina et al., 2008), a linear correlation between ρ^{ind} [-] and original grey-level (*GL*) is obtained by:

- subtracting the *GL*⁰ threshold value from the original *GL* values.
- **312** setting to zero all *GL* values lower than *GL*₀.
- 313 renormalizing the modified *GL* field.

314
$$\rho^{ind} = \begin{cases} GL > GL_0, & \frac{GL - GL_0}{\langle GL - GL_0 \rangle} \\ GL \le GL_0, & 0 \end{cases}$$
(13)

where notation $\langle \rangle$ indicates the mean value. GL_0 is adjusted based on the locations of pores that can be reliably identified on the image. By using *Fiji* (https://fiji.sc/), an open source 317 Java-based image processing package, it was possible to estimate the value of GL_0 for all three 318 materials.

In *Figure 4* are presented eight horizontal ρ^{ind} fields of the GW substrate extracted from the full-three dimensional scanned image. The resolution of the presented planes is 1024x1024 pixels ($\lambda_n = 2^{10}$), and they are equally distant in vertical direction (1.6 cm between two consecutive images). In Figure 5A are presented four horizontal ρ^{ind} fields of La Herreria soil – Horizon A, while in Figure 5B are presented horizontal fields of Horizon A20 (all images are 512x512 pixels - $\lambda_n = 2^9$).

325 4.2. UM model vs. Experimental data

The two-dimensional fields can be analyzed instead of the full-three dimensional one only if the statistical isotropy within the soil specimen is secured, which is the case here (explained further in the text). Note that $\rho_{s,min}^{ind}$ and $d_{g,min}$ are physical properties of the particular material and they are considered as unique for the whole sample (and every two-dimensional slice within it), while α and C_1 are statistical parameters determined for every horizontal plane individually.

332 4.2.1. UM model parameters

In case of the GW substrate, value of $\rho_{s,min}^{ind}$ is computed as $\rho_{s,min}^{ind} = \frac{\rho_{s,min}}{\rho_{bulk}} = \frac{2.2}{1.42} = 1.55$, where $\rho_{s,min}$ and ρ_{bulk} are experimentally determined, while in case of La Herreria soils $\rho_{s,min}^{ind}$ values are adjusted to fit measured percentages of sand, silt and clay particles (explained later in the text) because the corresponding $\rho_{s,min}$ values are missing (only ρ_{bulk} values are given). Clearly, due to its physical basis, value of $\rho_{s,min}^{ind}$ should not change significantly regardless of the soil type (adjusted values are $\rho_{s,min}^{ind} = 1.73$ and 1.54 for Horizon A and A20, respectively).

Value of $d_{g,min}$ depends on the range of scales that is analyzed. If focusing on the range of 340 341 scales covered by the scanned image, $d_{g,min}$ corresponds to the pixel size at λ_n (= 1024 and 512 for GW substrate and La Herreria soils, respectively), and thus it is equal to $d_{g,min} = L / \lambda_n \approx 60$ 342 / 1024 \approx 50 μm in case of the GW substrate, and $d_{g,min}$ = 16 / 512 \approx 32 μm for La Herreria 343 soils. For the full-range of scales, $d_{g,min}$ represents the minimal grain size that needs to be 344 approximately estimated if not measured. For the GW substrate, $d_{g,min} = 1 \ \mu m$ is adopted 345 based on the GW experimental data, while lower values of $d_{g,min} = 0.1 \ \mu m$ and 0.05 μm are 346 347 adopted for Horizons A and A20 of La Herreria soil, respectively, having on mind significant 348 percentages of clay particles in both cases (10% and 15%, respectively).

Finally, the UM parameters are determined by performing TM analysis on every $\rho^{ind}(\lambda)$ field 349 presented in Figure 4 and Figure 5. In Figure 6 is presented $log(\langle (\rho^{ind}(\lambda))^p \rangle)$ versus $log(\lambda)$ 350 for four different p values (0.2, 1.5, 2 and 3) and eight horizontal $\rho^{ind}(\lambda)$ fields of the GW 351 substrate presented in Figure 4. In all cases, scaling of $log(\langle (\rho^{ind}(\lambda))^p \rangle)$ (different symbols) 352 353 can be reasonably well interpreted with linear regressions. This is also the case with La 354 Herreria soils whose moment scaling behavior is illustrated in Figure 7A for Horizon A, and in Figure 7B for Horizon A20. The same kind of analysis is applied on the eight vertically 355 oriented fields ($\lambda_n = 2^{10}$) of the GW substrate (equal horizontal distance between the two 356 consecutive images). The quality of scaling is slightly better for vertical planes, but they are 357 358 not presented in the Figure. Vertical planes for La Herreria soils were not analyzed since the maximal resolution of the vertically oriented image is 256x256 pixels ($\lambda_n = 2^8$), which is 359 360 regarded as insufficient.

Based on the slopes of the obtained linear regressions, in Figure 8 are presented K(p)361 functions for all analyzed horizontal planes (different solid lines) of three different soils, 362 together with K(p) functions related to the vertical planes of the GW substrate (dashed lines in 363 364 Figure 8A). Since vertical and horizontal K(p) functions are overlapping in case of the GW 365 substrate, indicating similar values of α and C_1 (see **Table 1**), it is reasonable to assume the statistical isotropy within the GW specimen. Even though vertical images are not analyzed, 366 367 the same assumption is adopted for two remaining soils, having on mind the obtained horizontal K(p) functions are rather similar (values of α and C_1 are presented in **Table 2**). 368

Low values of C_1 (order of magnitude 10^{-2}) obtained for all three soils indicate the narrow range of ρ^{ind} values, meaning that ρ cannot be significantly larger than ρ_{bulk} (should be the case regardless of the soil type). Also, higher α values (closer to 2) point out more significant fluctuations of ρ around ρ_{bulk} , indicating the presence of different grain sizes with different densities.

374 4.2.2. Comparison with Experimental data

After determining values of the four model parameters, Equation (9) is firstly tested on the 375 ρ^{ind} field of the Hor. plane 4 of the GW substrate (see Figure 4). Since focusing on the range 376 of scales covered by the image, $d_{g,min} = 50 \ \mu m$ is used. Figure 9a-top illustrates comparison 377 378 between Equation (9) (UM model - solid line), Equation (4) (connected dots) that uses the counted number of $\rho^{ind}(\lambda) \geq \rho^{ind}_{s,min}$ values at different λ , and the truncated experimental 379 GSD data of the GW substrate (triangles). In Figure 9a-bottom are compared probability 380 381 density functions coming from the UM model (Equation 10 – solid lines) and measurements 382 (triangles). Good agreement between different analytical curves and truncated experimental data that consider only $d_g \ge 50 \ \mu m$ is obtained. 383

In Figure 9b-top is presented comparison between the same truncated experimental data (triangles) and Equation (9) applied on every ρ^{ind} field of Figure 4 (different solid lines), while in Figure 9b-bottom are illustrated the corresponding probability density functions (Equation 10).

Finally, in Figure 9c-top and Figure 9c-bottom the full range experimental data (squares) are compared with Equations (9) and (10), respectively, by using the same parameter values as in Figure 9b, with only difference that $d_{g,min} = 1 \ \mu m$ is adopted. The agreement between the Equation (9) / Equation (10) and both truncated and full-range experimental data is considered as satisfactory. The obtained family of curves creates reasonably narrow confidential zone around experimental points, verifying that way the analytical model proposed.

394 In Figure 10 is presented the same kind of analysis as in Figure 9, but for Horizons A and A20 of La Herreria soil. In Figure 10a and b are illustrated results obtained from ρ^{ind} fields of 395 Figure 5a (Horizon A) and Figure 5b (Horizon A20), respectively, where $d_{g,min} = 32 \ \mu m$ is 396 397 adopted, while in Figure 10c are compared full-range results for Horizon A (solid lines) and Horizon A20 (dashed lines) using $d_{g,min} = 0.1 \ \mu m$ and $d_{g,min} = 0.05 \ \mu m$, respectively. As 398 mentioned earlier in the text, values of $\rho_{s,min}^{ind}$ (1.73 and 1.54 for Horizons A and A20, 399 400 respectively) are adjusted so that the average percentages of sand, silt and clay particles, 401 computed based on the model results from Figure 10c, fit well with measured values. For 402 Horizon A, the average computed values of sand, silt and clay particles are 64%, 27% and 8% (about 1% of particles > 2 mm), respectively, while for Horizon A20 those values are 54%, 403 404 32% and 14%, respectively. The results obtained with a model are rather close to 405 measurements, confirming the UM model is valid.

406 The proposed UM approach was successfully evaluated on soils that cover wide range of407 grain sizes with significant percentages of both coarse and fine particles. To see limitations of

the model, it was additionally tested on a material with a rather uniform GSD curve, the Hostun sand HN31 extracted in the Drôme region in France (Bruchone et al, 2013). This sand is made of about 98% of quartz (grain densities are uniform), and it covers rather narrow spectrum of grain sizes ($0.2 \div 0.8$ mm). Preliminary results showed the proposed approach is not applicable on such material, having on mind the scaling of statistical moments in log-log scale significantly deviates from linear regression, and hence parameters α and C_1 cannot be determined.

415 4.3. UM model vs. PSF model

The PSF approach, a three-phase fractal-based GSD model firstly introduced by Perrier et al.
(1999), is also used for interpreting the experimental GSD curve of the GW substrate.
According to this model, the GSD can be described using the following expression (Bird et al.
2000):

420
$$P(d < d_g) = \left(\frac{d_g}{d_{g,max}}\right)^{3-D_{f,PSF}}$$
(14)

421 where $D_{f,PSF}$ is the fractal coefficient [-] whose optimal value can be determined from the 422 slope of the best fitting linear regression in logarithmic scale that goes through the 423 experimental GSD data and reaches 100 % at $d_{g,max}$. By adopting $d_{g,max} = 18$ mm in case of the 424 GW substrate, the optimal value of $D_{f,PSF} = 2.57$ is obtained (dashed line in Figure 8a).

To better illustrate the difference between the PSF (Equation 14) and the UM model (Equation 9), probability functions are presented in a log-log scale (Figure 11a). For the same values of *L*, $d_{g,min}$ and $\rho_{s,min}^{ind}$ as in Figure 9c, the best agreement between Equation (9) and measurements is obtained for $\alpha = 1.60$ and $C_I = 2.25 \times 10^{-2}$. These values are rather close to those describing the horizontal plane 4 of the GW substrate (see Table 1), confirming the physical basis of the proposed approach.

To test whether the UM model can be used for soils with multi-modal GSD curves that are 431 described with multiple fractal dimensions, Equation (9) is compared with experimental data 432 433 of Walla Walla (WW) soil, described in Bittelli et al. (1999) by means of three different 434 fractal dimensions related to distributions of sandy, silty and clayey particles. According to the U.S. soil taxonomy, this soil is Typic Haploxeroll with 8.3% of sand, 78.4% of silt and 435 436 13.3% of clay. Due to the lack of scanned soil images and other relevant information, the UM 437 model parameters could not be determined as explained earlier in the text, but only roughly estimated / adjusted. Figure 11b shows that for $\rho_{s,min}^{ind} = 1.72$, $\alpha = 1.05$, $C_1 = 4.5 \times 10^{-2}$, and L / 438 $d_{g,min} = 1 \text{ mm} / 1 \times 10^{-4} \text{ mm}$ (from the graph), Equation (9) provides rather good agreement 439 with WW experimental points. Since the relevant information are missing, these results 440 441 should be taken with a grain of salt and understood only as a test if the model is capable of 442 interpreting multi-modal GSD curves.

443

444 Compared to Equation (14) which considers the fixed fractal dimension $D_{f,PSF}$, Equation (9) 445 takes into account the fractal dimension that changes with d_g (the co-dimension function). Therefore, for the UM approach it is quite important to know the total range of scales 446 investigated (L / $d_{g,min}$). Even though the PSF model (Equation 14) is more convenient for 447 practical application, since it uses only two parameters that can be determined quite easily, the 448 449 UM model proposed in this work (Equations 9 and 10) shows better agreement with experimental GSDs for soils with wide spectrum of grain sizes whose distribution cannot be 450 451 accurately described with a simple fractal (power) law.

452 **5.** Conclusion

This work shows that the up-scaling approach presented can be used for predicting the GSD of a certain material based on its scanned micro-structure which represents a density indicator

field. By reducing the resolution of the scanned image, density indicator values above the 455 456 fixed threshold are treated at each resolution as a cumulative representation of solid particles 457 of diameter equal to or larger than the corresponding pixel size. The quantity of values above 458 the threshold at different resolutions can be analytically described by means of the Universal 459 Multifractals (UM), leading to the new multifractal-based GSD model. The model uses four 460 parameters, where two of them are physical (the minimal grain diameter and the ratio between the minimal grain density and dry bulk density), while the other two are statistical UM 461 462 parameters that characterize the spatial heterogeneity of the soil density field.

463 An innovative approach proposed in this work was tested on three different materials containing significant percentages of both coarse and fine particles: an unconventional 464 volcanic granular material used for covering green roofs, and two horizons of La Herreria soil 465 466 collected from the low land mountain area of Sierra de Guadarrama in Spain. By performing 467 the Trace Moment analysis on two-dimensional scanned soil images (density indicator field) 468 of different materials, the statistical parameters (C_1 and α) were determined and families of 469 analytical GSD curves were obtained for each material, showing a good agreement with 470 experimental data. Additionally, the UM approach was tested on Houston quartz sand and 471 results showed that such a uniform GSD cannot be reliably estimated by means of the presented methodology. In all cases, the analysis was applied on two-dimensional images to 472 473 save computational time and memory, which is legit if the statistical isotropy within the 474 specimen is secured.

475 Results showed the values of C_1 are the same order of magnitude for different soils (10⁻²) 476 because density of an individual grain cannot be significantly larger than the dry bulk density, 477 regardless of the soil type. Also, it was proved that lower C_1 and higher α are related to GSD 478 curves that change gradually with respect to grain sizes, while lower α and higher C_1 are 479 related to the steeper GSD curve in its central part and stronger curvature close to the break480 onto the finer particles.

Finally, the proposed four parameters UM model was compared with the fractal-based two parameters PSF model, showing better agreement with the Green Wave experimental data. Its advantage was additionally emphasized by showing excellent agreement with multi-modal GSD curve of Walla Walla soil that is described in the literature by means of three different fractal dimensions related to sandy, silty and clayey particles.

486

487 Appendix: Analogy between the up-scaling approach and the dry sieving

- 488 method
- The analogy between Equation (4) (the up-scaling approach) and the dry sieving method canbe derived under the following assumptions:

491 1. Three-dimensional space is considered (E = 3) since the dry sieving method is based on 492 the grain masses.

- 493 2. All grains have the same shape and density ρ_s .
- 494 3. A discrete number of sieves is used, and hence the size of voids on each sieve follows the 495 size of voxels $l(\lambda_{n-i})$ at different λ_{n-i} , where $i = [0 \div n]$.

The total mass of the specimen (M_{total}) is placed on the cascade of sieves arranged in a descending order (the largest void size is on the top, while the smallest one is on the bottom). If all grains are distributed on the corresponding sieves $i = [0 \div n]$, it can be assumed that those staying on a certain sieve have diameter equal to the size of sieve voids. Therefore, the following can be written for i = 0:

501
$$M(d \ge l(\lambda_n)) = M_{total} = C_V \rho_s \sum_{j=0}^n N_g(\lambda_{n-j}) l(\lambda_{n-j})^3$$
(A1)

where C_V is the volume shape coefficient [-] (for cube $C_V = 1$, for sphere $C_V = \pi / 6$), $N_g(\lambda_{n-j})$ is the number of grains [-] that stay on the sieve of void size $l(\lambda_{n-j})$ [L] $(j = 0 \div n)$. Since grains that stay on the sieve $l(\lambda_n)$ are assumed to have identical diameters, $l(\lambda_n)^3$ can be pulled outside the sum, and the following is obtained:

506
$$M(d \ge l(\lambda_n)) = C_V \rho_s l(\lambda_n)^3 \sum_{j=0}^n N_g(\lambda_{n-j}) \left(\frac{l(\lambda_{n-j})}{l(\lambda_n)}\right)^3$$
(A2)

507
$$N_g^{cum}(\lambda_n) \approx C_V \sum_{j=0}^n N_g(\lambda_{n-j}) \left(\frac{l(\lambda_{n-j})}{l(\lambda_n)}\right)^3$$
 (A3)

where $N_g^{cum}(\lambda_n)$ is the cumulative number of grains equal to or larger than $d_{g,min} = l(\lambda_n)$. Following Equation (A2), the cumulative mass of all grains equal to or larger than $l(\lambda_{n-i})$ can be expressed as the following:

511
$$M(d \ge l(\lambda_{n-i})) = C_V \rho_s l(\lambda_{n-i})^3 \sum_{j=i}^n N_g(\lambda_{n-j}) \left(\frac{l(\lambda_{n-j})}{l(\lambda_{n-i})}\right)^3$$
(A4)

512
$$N_g^{cum}(\lambda_{n-i}) \approx C_V \sum_{j=i}^n N_g(\lambda_{n-j}) \left(\frac{l(\lambda_{n-j})}{l(\lambda_{n-i})}\right)^3$$
 (A5)

513 By introducing Equation (A3) into (A2) and Equation (A5) into (A4), and having on mind 514 that $d_g = l(\lambda_{n-i})$, the GSD can be expressed as:

515
$$P(d < d_g) = 1 - \frac{M(d \ge l(\lambda_{n-i}))}{M(d \ge l(\lambda_n))}$$
(A6)

516
$$P(d < d_g) = 1 - \frac{N_g^{cum}(\lambda_{n-i})}{N_g^{cum}(\lambda_n)} \left(\frac{l(\lambda_{n-i})}{l(\lambda_n)}\right)^3$$
(A7)

517 Note that $N_g^{cum}(\lambda_{n-i})$ and $N_g^{cum}(\lambda_n)$ correspond to $N(\rho^{ind}(\lambda) \ge \rho_{s,min}^{ind})$ and 518 $N(\rho^{ind}(\lambda_n) \ge \rho_{s,min}^{ind})$ in Equation (4), respectively, while $\frac{l(\lambda_{n-i})}{l(\lambda_n)} = \frac{\lambda_n}{\lambda_{n-i}}$. Therefore, 519 Equation (A7) is identical to Equation (4) for E = 3, just in a discrete form.

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524

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Figure 1. Change of the two-dimensional $\rho^{ind}(\lambda)$ field with λ at: (A) $\lambda = 128$; (B) $\lambda = 64$; (C) λ = 32; (D) $\lambda = 16$; (E) $\lambda = 8$; (F) $\lambda = 4$. A turquoise color platform signifies a fixed threshold $\rho_{s,min}^{ind} = 1.55$, while the one with transparent color signifies a resolution dependent threshold $\lambda^{\gamma}(\gamma = 0.211)$



Figure 2. A) Equation (1) (dashed lines) calculated for: $C_I = 2.23 \times 10^{-2}$, $\alpha = 1.67$ (values that characterize $\rho^{ind}(\lambda)$ field in Figure 1) and different values of γ (including $\gamma = 0.211$); B) Equation (5) (solid line) calculated for the same values of C_I and α and the fixed value of $\rho_{s,min}^{ind} = 1.55$.



671

Figure 3. Behavior of the proposed GSD model when changing values of: a) C_1 ; b) α ; c) $\rho_{s,min}^{ind}$; d) $d_{g,min}$. Initial parameter values (solid line in each graph) are $C_1 = 1.85 \times 10^{-2}$, $\alpha = 1.3$, $\rho_{s,min}^{ind} = 1.55$ and $d_{g,min} = 1 \times 10^{-3}$ mm

676



Figure 4. Eight two-dimensional horizontal ρ^{ind} fields (1024 x 1024 pixels), extracted from the original three-dimensional grey scale image, that are equally distant along the specimen height

(A) La Herreria - Horizon A

A - Hor. plane 1 A - Hor. plane 2 A - Hor. plane 3 A - Hor. plane 4 -7 -6 -5 -4 -3 -2 -6 -5 -4 -3 -2 -4 -3 -2 2





699 Figure 5. (A) Four equally distant horizontal ρ^{ind} fields (512 x 512 pixels) of Horizon A of La

700 Herreria soil, extracted from the original three-dimensional grey scale image; (B) same as in

701 (A) just for Horizon A20



Figure 6. Scaling of statistical moments of eight horizontal $\rho^{ind}(\lambda)$ fields of the GW substrate presented in Figure 4

(A) La Herreria - Horizon A





Figure 7. (A) Scaling of statistical moments of four horizontal $\rho^{ind}(\lambda)$ fields of La Herreria soil

-9-0-0-0

 $log(\lambda)$

0.00

Ő

-0 •• 0-

 $log(\lambda)$

0.00

 $log(\lambda)$

-9-0-0-0

0.00

-9-0-0-0

(Horizon A) presented in Figure 5a; (B) same as in (A) just for Horizon A20 (fields presented

in Figure 5b)

0.00

Ō

 $log(\lambda)$



Figure 8. Moment scaling functions *K(p)* obtained by applying TM technique on: (A) eight
horizontal fields from Figure 4 (different solid lines), and eight vertically oriented fields
(dashed lines); (B) four horizontal fields from Figure 5a; (C) four horizontal fields from
Figure 5b



Figure 9. Comparison between: A-top) Equation (9) applied on the GW Hor. plane 4 - solid line (L = 60 mm, $\rho_{s,min}^{ind}$ = 1.55, $d_{g,min}$ = 50 µm, C_I = 2.23x10⁻², α = 1.67), Equation (4) applied on the same field – connected dots ($\rho_{s,min}^{ind} = 1.55$, $\lambda_n = 1024$), and truncated experimental GSD data ($d_g \ge 50 \ \mu m$) - triangles; A-bottom) Equation (10) computed with the same parameter values as in A-top, and truncated experimental GSD data (triangles); B-top) Equation (9) applied on eight $\rho^{ind}(\lambda)$ fields from Figure 4 – solid lines (L, $\rho^{ind}_{s,min}$ and $d_{s,min}$ identical as in A, α and C_1 presented in **Table 1**), and truncated experimental GSD data (triangles); B-bottom) Same as in A-bottom just for all fields from Figure 4; C-top) Same as in B-top just for $d_{g,min} = 1 \ \mu m$ (solid lines), and full-range experimental GSD data ($d_g \ge 1 \ \mu m$) - squares; C-bottom) Same as in B-bottom just for $d_{g,min} = 1 \mu m$ (solid lines), and full-range experimental GSD data ($d_g \ge 1 \ \mu m$) – squares



Figure 10. Comparison between: (A) Equation (9) applied on four Horizon A fields from Figure 5A (L = 16 mm, $\rho_{s,min}^{ind} = 1.73$, $d_{g,min} = 32 \text{ µm}$, α and C_1 from Table 2), and Equation (4) (box-counting method) applied on same fields ($\lambda_n = 512 \approx L/d_{g,min}$); (B) Same as in (A) just for Horizon A20 ($\rho_{s,min}^{ind} = 1.54$); (C) Equation (9) applied on four fields from Figure 5A (parameters same as in (A), only $d_{g,min} = 0.1 \text{ µm} - \text{solid lines}$) and on four fields from Figure 5B (parameters same as in (A), only $d_{g,min} = 0.05 \text{ µm} - \text{dashed lines}$)

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Figure 11. (A) Comparison between the experimental GSD data of the GW substrate (squares) and Equations (9) (solid lines) and (14) (dashed lines) computed with the adjusted parameter values that fit the measurements ($C_1 = 2.25 \times 10^{-2}$ and $\alpha = 1.60$ for UM model, and $D_{f,PSF} =$ 2.57 for PSF model); (B) same as in (A) just for Walla Walla soil taken from Bittelli et al. (1999) – L = 1 mm, $d_{g,min} = 1 \times 10^{-4}$ mm, $\rho_{s,min}^{ind} = 1.72$, $\alpha = 1.05$, $C_1 = 4.5 \times 10^{-2}$ (PSF model is not illustrated since it clearly deviates from the measurements)

Table 1. Determined UM parameters for eight $\rho^{ind}(\lambda)$ fields of the GW substrate presented in

832 Figure 4

	Horizontal plane								
[1	2	3	4	5	6	7	8	
<i>C</i> ₁	9.34E-03	9.93E-03	1.64E-02	2.23E-02	1.66E-02	2.72E-02	1.45E-02	1.93E-02	
α	1.93	1.96	1.83	1.67	1.80	1.61	1.85	1.66	
-	Vertical plane								
	1	2	3	4	5	6	7	8	
\mathcal{I}_{1}	2.66-02	2.21E-02	1.74E-02	1.72E-02	1.67E-02	1.63E-02	2.06E-02	2.37E-02	
α	1.56	1.72	1.75	1.75	1.76	1.82	1.78	1.60	

Table 2. Determined UM parameters for eight $\rho^{ind}(\lambda)$ fields of La Herreria soil (Horizons A

and A20) presented in Figure 5

	Horizon A								
	1	2	3	4					
C_1	3.36E-02	3.55E-02	3.03E-02	2.60E-02					
α	1.15	1.25	1.30	1.37					
	Horizon A20								
	1	2	3	4					
C_1	1.62-02	1.63E-02	1.73E-02	1.62E-02					
α	1.46	1.40	1.48	1.44					